

## Counting statistics of single-electron capture by a dynamic quantum dot

L. Fricke<sup>1</sup>, M. Wulf<sup>1</sup>, B. Kaestner<sup>1</sup>, V. Kashcheyevs<sup>2</sup>, J. Timoshenko<sup>2</sup>, P. Nazarov<sup>2</sup>, F. Hohls<sup>1</sup>, P. Mirovsky<sup>1</sup>, B. Mackrodt<sup>1</sup>, R. Dolata<sup>1</sup>, T. Weimann<sup>1</sup>, K. Pierz<sup>1</sup>, H.W. Schumacher<sup>1</sup>

<sup>1</sup>*Physikalisch-Technische Bundesanstalt (PTB), Bundesallee 100, D-38116 Braunschweig, Germany*

<sup>2</sup>*Faculty of Physics and Mathematics, University of Latvia, Riga LV-1002, Latvia*

A renewed International System of Units (SI) strongly demands for a quantum-based current source relating the output current to the elementary charge  $e$  whose value will be fixed at redefinition [1]. A highly promising candidate for such a current source is the non-adiabatic single-electron pump [2, 3] exploiting a dynamic quantum dot forming out of a two-dimensional electron gas. Due to the dynamic tunnel barriers between the dot and the source/drain leads the dot can be driven by high frequencies since its population is not limited by tunneling constants. Recent measurements have qualified the pump accuracy to be better than  $10^{-6}$  [4]. However, the exact initialization mechanism of these dynamic dots is still subject of current research.

Recently, a theoretical model of the probability of charge capture has been developed [5], including mainly three relevant energy scales for this process. These are the temperature ( $kT$ ), the finite time scale for suppression of backtunneling (expressed by  $\Gamma_c$ ) as well as the coupling of the rising source barrier to the energy levels of the quantum dot due to electrostatic cross talk ( $\Delta_{ptb}$ ). In the limit  $\Gamma_c, \Delta_{ptb} \rightarrow 0$  the probability of charge capture follows a thermal distribution. In a second limit with  $kT, \Gamma_c \rightarrow 0$ , the previously predicted decay-cascade model [6] is reproduced.

To investigate these regimes and to identify the relevant processes in our samples, we combine such a dynamic quantum dot with highly-sensitive electrometers and perform counting measurements on the number of charges initialized on the dot and subsequently transferred to a measurement node. Using this architecture we are able to distinguish between these two limits and identify the decay-cascade regime as the dominating mechanism of charge capture in our sample [7]. Additionally, based on the relevant mechanism, we propose different strategies for further improvement.

Furthermore, an overview about the actual status of the self-referenced current source including an error-accounting scheme [8] will be given.

- [1] 24th Resolution of the CGPM, available online via <http://www.bipm.org/utls/common/pdf/24.CGPM.Resolutions.pdf>.
- [2] M. Blumenthal et al., *Nature Physics* **3**, 343 (2007).
- [3] B. Kaestner et al., *Phys. Rev. B* **77**, 153301 (2008).
- [4] S. Giblin et al., *Nature Communications* **3**, 930 (2012).
- [5] V. Kashcheyevs, J. Timoshenko, *Phys. Rev. Lett.* **109**, 216801 (2012).
- [6] V. Kashcheyevs, B. Kaestner, *Phys. Rev. Lett.* **104**, 186805 (2010).
- [7] L. Fricke et al., *Phys. Rev. Lett.* accepted and arXiv:1211.1781.
- [8] M. Wulf, *Phys. Rev. B* **87**, 035312 (2013).