

## Transmission through silicene quantum barriers

Diana M. M. Gustin<sup>1</sup>, Marcos R.S. Tavares<sup>1</sup>, G.-Q. Hai<sup>2</sup>, and P. Vasilopoulos<sup>3</sup>

<sup>1</sup>Centro de Ciências Naturais e Humanas, Univ. Federal do ABC, 09210-170, S. André, SP, Brazil

<sup>2</sup>Instituto de Física de São Carlos, Universidade de São Paulo, 13560-970, São Carlos, SP, Brazil

<sup>3</sup>Department of Physics, Concordia University, 7141 Sherbrooke Ouest, Montreal, Canada H4B 1R6

Silicene is a monolayer of silicon atoms that form a two-dimensional honeycomb lattice similar to graphene's. It has attracted considerable attention due to its exotic electronic structure and its compatibility with current silicon-based electronic technology. The low-energy physics of silicene is described by Dirac electrons with relative large spin-orbit interaction,  $\lambda = 3.9 \text{ meV}$ , due to its buckled structure<sup>1</sup> with  $l = 2.3 \times 10^{-2} \text{ nm}$  the half distance between the A and B sublattices. Its band structure can be controlled externally by an electric field  $E_z = 300 \text{ meV/nm}$ . We theoretically explore resonant features in the electronic transmission  $T$  through barriers in silicene by studying  $T$  as a function of the electron's energy  $E$  (meV) and its angle of incidence  $\theta$ . The incident electrons are assumed spin polarized and the barriers are formed by contact potentials. Our results show that the applied electric field can result in spin-resolved transmission channels. In addition, we critically compare the results with those through barriers in graphene. Important differences from the graphene's case<sup>2</sup> show up and are highlighted in the transmission  $T^\uparrow$  and polarization  $P = |T^\uparrow - T^\downarrow|/|T^\uparrow + T^\downarrow|$ . A  $(E, \theta)$  contour plot of  $T$ , for incident *spin-up electrons*, through a silicene barrier of length  $L=110 \text{ nm}$  and height  $V=100 \text{ meV}$ , is shown in Fig. 1(a). Notice the periodicity along the  $E$  axis for large  $\theta$ . Figure 1(b) is the same contour as in Fig. 1(a) but for the values  $\lambda = 0 \text{ meV}$  and  $E_z = 0$  which pertain to a barrier on suspended graphene. Notice the absence of the transmission channel below 20 meV shown in Fig. 1(a). The contour plot in Fig. 1(c) shows the spin polarization  $P$  for a barrier on silicene.  $P$  increases strongly for energies near the top of the barrier and also for  $E_+ = (\lambda + l * E_z)$  and  $E_- = -(\lambda - l * E_z)$ .--- This work is supported by CAPES, CNPq, and FAPESP, Brazil, and the NSERC grant OGP012756, Canada.

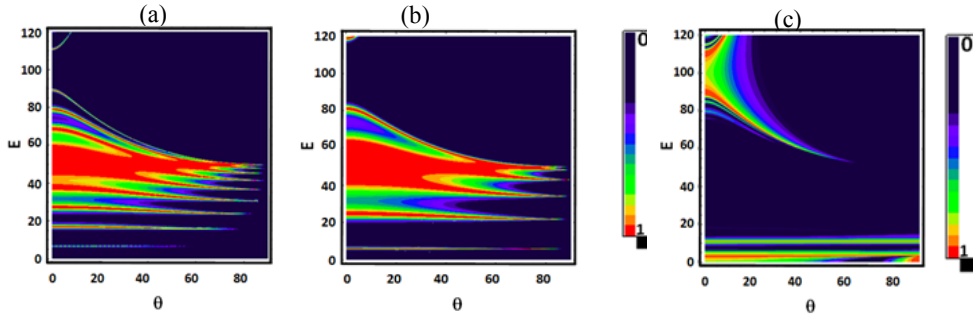


Figure 1. (a)  $(E, \theta)$  contour plot of the transmission through a silicene barrier. (b) As in (a) for a barrier on suspended graphene. (c)  $(E, \theta)$  contour plot of the polarization for a silicene barrier.

[1] Motohiko Ezawa, Phys. Rev. Lett. **109**, 055502 (2012).

[2] J. Milton Pereira Jr. *et al.*, Appl. Phys. Lett. **90**, 132122 (2007).